

Enakomerna kompresija, kompresijski modul $\kappa = \frac{E}{3(1-2\nu)}$.

Nestisljivi material $\nu = \frac{1}{2}$, vrednosti Poissonovega količnika $\nu \in (-1, \frac{1}{2}]$.

Termoelastičnost

$$\underline{\underline{\epsilon}} = \underline{\underline{\epsilon}}_T + \underline{\underline{\epsilon}}_E$$

Termalni raztezek $\underline{\underline{\epsilon}}_T = \alpha \Delta T \underline{\underline{I}}$.



$$\underline{\underline{\epsilon}} = \underbrace{\frac{1+\nu}{E} \underline{\underline{t}} - \frac{\nu}{E} (\text{sl } \underline{\underline{t}}) \underline{\underline{I}}}_{\underline{\underline{\epsilon}}_E} + \alpha \Delta T \underline{\underline{I}}$$

Primer: izračun napetostnega stanja elastičnega vključka v togi matriki.

$$\underline{\underline{\epsilon}} = \underline{\underline{\epsilon}}_E + \underline{\underline{\epsilon}}_T$$

$$\underline{\underline{\epsilon}} = 0$$

$$\underline{\underline{\epsilon}}_E = -\underline{\underline{\epsilon}}_T = -\alpha \Delta T \underline{\underline{I}}$$

$$\frac{1+\nu}{E} \underline{\underline{\tau}} - \frac{\nu}{E} \text{sl } \underline{\underline{\tau}} \underline{\underline{I}} = -\alpha \Delta T \underline{\underline{I}}$$

$\underline{\underline{\tau}}$ je sferični tenzor $\underline{\underline{\tau}} = p \underline{\underline{I}}$

Napetostno stanje je hidrostatično.

$$\text{sl} \left(\frac{1+\nu}{E} \underline{\underline{\tau}} - \frac{\nu}{E} \text{sl } \underline{\underline{\tau}} \underline{\underline{I}} \right) = \text{sl} (-\alpha \Delta T \underline{\underline{I}})$$

$$\frac{1+\nu}{E} \text{sl } \underline{\underline{\tau}} - \frac{\nu}{E} \text{sl } \underline{\underline{\tau}} \cdot 3 = -\alpha \Delta T \cdot 3$$

3

$$\left(\frac{1+\nu}{E} - 3 \frac{\nu}{E} \right) \text{sl } \underline{\underline{\tau}} = -3 \alpha \Delta T$$

$$\frac{1-2\nu}{E} \text{sl } \underline{\underline{\tau}} = -3 \alpha \Delta T \Rightarrow \text{sl } \underline{\underline{\tau}} = -\frac{3E}{1-2\nu} \alpha \Delta T$$

$$\underline{\underline{\sigma}} = - \frac{3E\alpha}{1-2\nu} \Delta T \underline{\underline{1}}$$

$$\underline{\underline{\epsilon}} = \frac{1+\nu}{E} \underline{\underline{\sigma}} - \frac{\nu}{E} \underline{\underline{\sigma}} \underline{\underline{1}}$$

Ravninska napetost

$$\sigma_{13} = \sigma_{23} = \sigma_{33} = 0$$

$$\epsilon_{11} = \frac{1}{E} \sigma_{11} - \frac{\nu}{E} \sigma_{22}$$

$$\epsilon_{22} = \frac{1}{E} \sigma_{22} - \frac{\nu}{E} \sigma_{11}$$

$$\epsilon_{13} = 0 ; \quad \epsilon_{23} = 0 ; \quad \epsilon_{12} = \frac{1+\nu}{E} \tau_{12}$$

$$\epsilon_{33} = - \frac{\nu}{E} \sigma_{11} - \frac{\nu}{E} \sigma_{22}$$

Ravninska deformacija

$$\epsilon_{13} = \epsilon_{23} = \epsilon_{33} = 0$$

$$\epsilon_{11} = \frac{1}{E} t_{11} - \frac{\nu}{E} t_{22} - \frac{\nu}{E} t_{33}$$

$$\epsilon_{22} = \frac{1}{E} t_{22} - \frac{\nu}{E} t_{11} - \frac{\nu}{E} t_{33}$$

$$0 = \epsilon_{33} = \frac{1}{E} t_{33} - \frac{\nu}{E} t_{11} - \frac{\nu}{E} t_{22} \Rightarrow t_{33} = \nu(t_{11} + t_{22})$$

$$\epsilon_{11} = \frac{1}{E} t_{11} - \frac{\nu}{E} t_{22} - \frac{\nu}{E} (\nu(t_{11} + t_{22})) = \frac{1-\nu^2}{E} t_{11} - \frac{\nu(1+\nu)}{E} t_{22}$$

$$\epsilon_{22} = \frac{1-\nu^2}{E} t_{22} - \frac{\nu(1+\nu)}{E} t_{11}$$

$$\frac{1-\nu^2}{E} = \frac{1}{E'} \quad \frac{\nu(1+\nu)}{E} = \frac{\nu'}{E'}$$

$$E' = \frac{E}{1-\nu^2}$$

$$\nu' = \frac{\nu(1+\nu)E'}{E} = \frac{\nu(1+\nu)E}{(1-\nu^2)E} = \frac{\nu}{1-\nu}$$

$$1-\nu^2 = (1-\nu)(1+\nu)$$

$$0 = \epsilon_{23} = \frac{1+\nu}{E} t_{23} \Rightarrow t_{23} = 0 ; \quad \epsilon_{13} = 0$$

$$\epsilon_{12} = \frac{1+\nu}{E} t_{12}$$

$$\frac{1+\nu}{E} = \frac{1+\nu'}{E'}$$

Modificirana modula

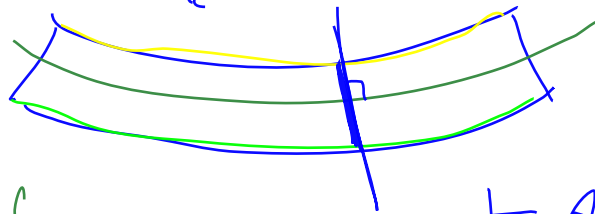
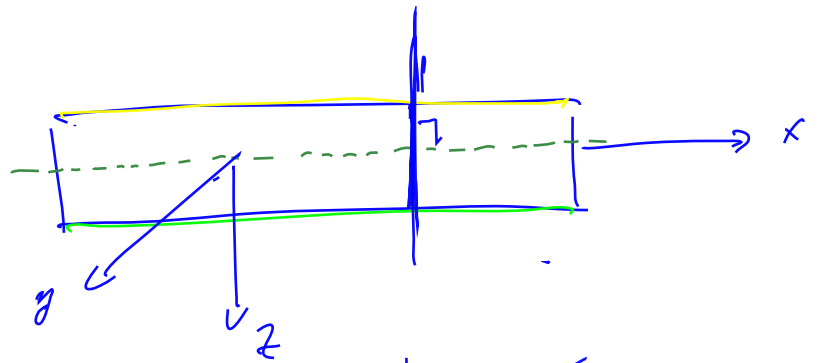
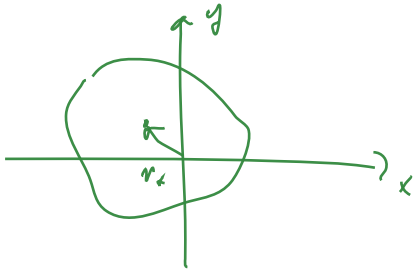
$$\hat{E} = \frac{E}{1-\nu^2}, \quad \hat{\nu} = \frac{\nu}{1-\nu}$$

$$\frac{1+\nu'}{E'} = \frac{(1 + \frac{\nu}{1-\nu})(1-\nu')}{E} = \frac{1}{E} \left(\frac{1-\nu^2 + \nu(1+\nu)}{1-\nu^2 + \nu + \nu^2} \right) = \frac{1+\nu}{E}$$

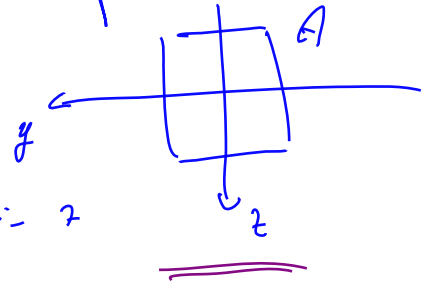
Ravninska deformacija je ekvivalentna ravninskoj napetosti s modificiranimi modula.

Upogib nosilca

Nevtralna os, centralna os.



$$\vec{r}_* = \frac{1}{A} \int_A \vec{r} dA$$



A simetričen glede na osi y i z

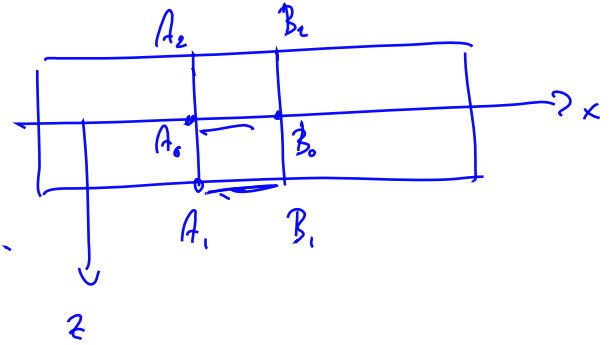
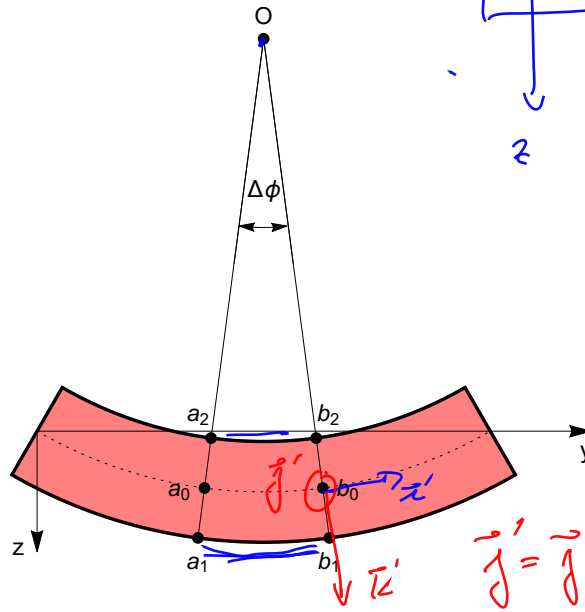
centralna os je os x = 1 $\vec{r}_ = (x, 0, 0)$*

$$z_* = \frac{1}{A} \int_A z dA = 0$$

Inženirska teorija nosilcev:

- Dolžina nosilca je bistveno večja od lateralnih dimenzij.
- Nosilec je simetričen glede na ravnini xz in xy .
- Za homogeni nosilec se nevtralna os ujema s centralno osjo.
- Ravnine pravokotne na nevtralno os v referenčni legi se deformirajo v ravnine, ki so pravokotne na deformirano nevtralno os.
- V ravninah, ki so pravokotne na nevtralno ni deformacij.

Deformacija vlaknen



$$|A_0 A_1| = z$$

$$|A_0 B_1| = -z$$

$$R = |O a_0| = |O b_0|$$

Slika 1: Upogib nosilca, nevtralna os je označena črtkano.

$$\epsilon = \frac{|a_1 b_1| - |A_1 B_1|}{|A_1 B_1|} = \frac{|a_1 b_1|}{|A_1 B_1|} - 1 = \frac{|a_1 b_1|}{|a_0 b_0|} - 1$$

$$|A_1 B_1| = |A_0 B_0| = |a_0 b_0| \quad |a_0 b_0| = R \phi$$

$$|O a_1| = |O a_2| = R + z$$

$$|a_1 b_1| = (R + z) \phi$$

$$\epsilon = \frac{(R + z) \phi}{R \phi} - 1 = \frac{z}{R}$$

$$\hat{\epsilon} = \frac{|a_1 b_1|}{|a_0 b_0|} - 1$$

$$\hat{\epsilon} = \frac{z}{R}$$

$$|O a_2| = R + z$$

$$\epsilon(x) = \frac{z}{R(x)}$$

Deformacija vlaknen $\epsilon = \frac{z}{R}$.

$$\sigma = E \epsilon = \frac{E}{R} z$$

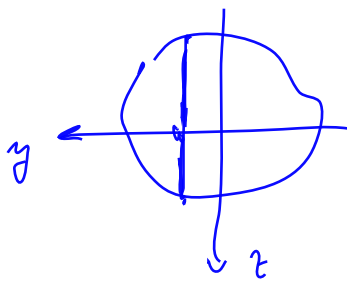
homogeni napetost

Napetost $\sigma = z \frac{E}{R}$.

$$\vec{\tau} = \sigma \cdot \vec{z}'$$

$$\vec{0} = \int_A \vec{\tau} dA = \int_A \sigma \vec{z}' dA = \int_A \frac{E}{R} z dA = \frac{E}{R} \int_A z dA = 0$$

$$\vec{M} = \int_A d\vec{M} = \int_A \vec{r} \times \vec{\tau} dA = \int_A (y\vec{j}' + z\vec{k}') \times \sigma \vec{z}' dA =$$



$$\vec{r} = y\vec{j}' + z\vec{k}'$$

$$= \frac{E}{R} \int_A (y z \vec{k}' + z^2 \vec{j}') dA = \left(\frac{E}{R} \int_A y z dA \right) \vec{k}' + \frac{E}{R} \int_A z^2 dA \vec{j}'$$

$\int_A y z dA = 0$ zaradi predpostavljene simetrije preseka

$$\vec{M} = \frac{E}{R} \int_A z^2 dA \vec{j}'$$

$$\vec{M} = \frac{E}{R} \int_A z^2 dA \vec{j}$$

$$M = \frac{E}{R} \int_A z^2 dA$$

Enaka Bernoullijevo lomnata.

$$\int_A z dA \text{ ploščinski moment druge vrste} = I_Y = I$$

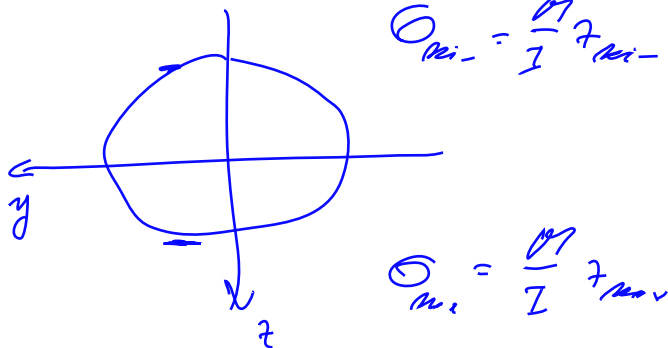
$$M = \frac{E I}{R}$$

$$\sigma = \frac{E}{R} z$$

$$\sigma = \frac{M}{I} z$$

Upogibni moment

Izračun upogibnega momenta.



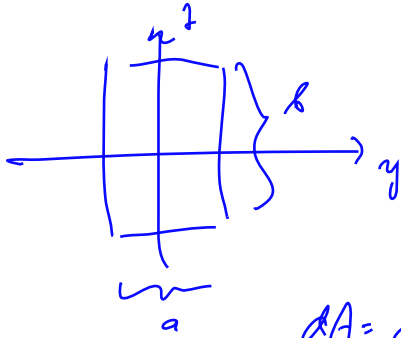
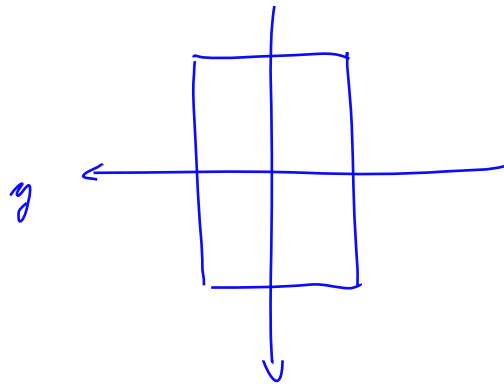
Euler - Bernoullijeva enačba $M = \frac{EI}{R}$.

Zveza $\sigma = E\epsilon = \frac{M}{I}z$.

$$\sigma = \frac{M}{I} z$$

Ploskovni moment

Ploskovni moment pravokotnika, $I = ab^3/12$.



$$I = \int_A z^2 dA = \int_{-a/2}^{a/2} dy \int_{-b/2}^{b/2} z^2 dz = 2a \int_0^{b/2} z^2 dz$$

$$dA = dy dz$$

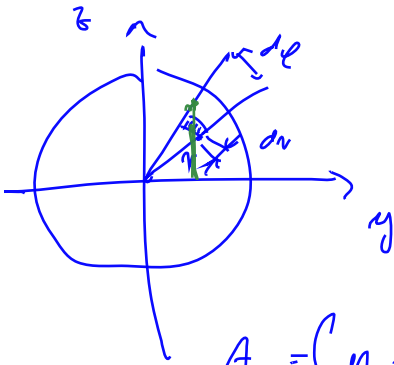
$$= 2a \left[\frac{1}{3} z^3 \right]_0^{b/2} = 2a \frac{1}{3} \cdot \frac{b^3}{8}$$

$$I = \frac{1}{12} a b^3$$

$$\int_{x_1}^{x_2} f(x) dx = F(x_2) - F(x_1)$$

$$F'(x) = f(x)$$

Ploskovni moment kroga $I = \pi R^4/4$.



$$I = \int_A z^2 dA \quad z = r \sin \varphi$$

$$dA = r dr d\varphi$$

$$A = \int_A dA = \int_0^R r dr \int_0^{2\pi} d\varphi = \left[\frac{1}{2} r^2 \right]_0^R \cdot 2\pi = \pi R^2$$

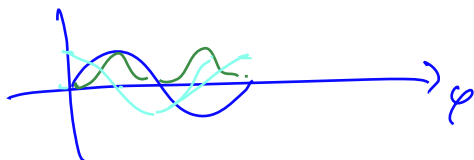
$$I = \int_A r^2 \sin^2 \varphi dA = \int_0^R r^3 dr \int_0^{2\pi} \sin^2 \varphi d\varphi = \left[\frac{1}{4} r^4 \right]_0^R \cdot \pi$$

$$\int_0^{2\pi} \sin^2 \varphi d\varphi = \int_0^{2\pi} \cos^2 \varphi d\varphi$$

$$2 \int_0^{2\pi} \sin^2 \varphi d\varphi = \int_0^{2\pi} \cos^2 \varphi d\varphi + \int_0^{2\pi} \sin^2 \varphi d\varphi$$

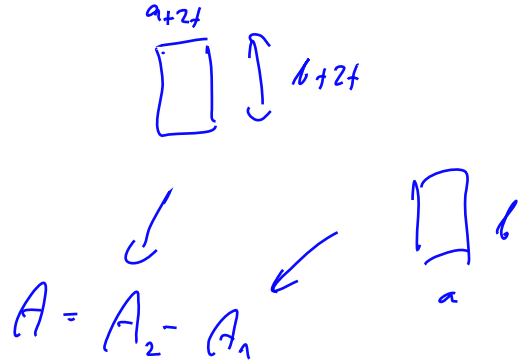
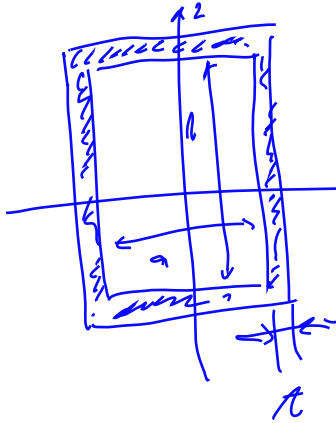
10

$$= \int_0^{2\pi} d\varphi = 2\pi \Rightarrow \int_0^{2\pi} \sin^2 \varphi d\varphi = \pi$$



$$I = \frac{1}{4} \pi R^4$$

Ploskovni momenti tankostenskih presekov.



$$I = \int_A z^2 dA = \int_{A_2} z^2 dA - \int_{A_1} z^2 dA =$$

$$\frac{1}{12} (a+2t)(b+2t)^3 - \frac{1}{12} ab^3$$

$$t \ll a; \quad t \ll b$$

$$I = \frac{1}{12} \left[(a+2t)(b^3 + 3b^2 \cdot 2t + 3 \cdot b \cdot (2t)^2 + (2t)^3) - ab^3 \right]$$

$$\frac{ab^3 + 2tb^3 + 3ab^2 \cdot 2t + 6b^2 t^2}{12}$$

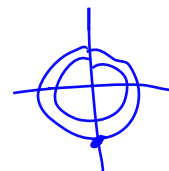
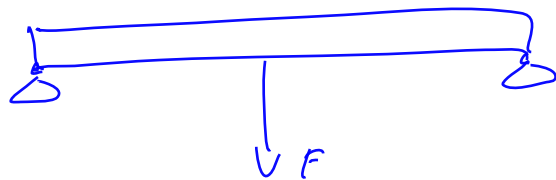
$$\approx \frac{1}{12} [2tb^3 + 6ab^2 t] = \frac{1}{6} (b^3 + 3ab^2) t \quad \checkmark$$



$$I = \int_{A_2} z^2 dA - \int_{A_1} z^2 dA = \frac{1}{4} \pi (R+t)^4 - \frac{1}{4} \pi R^4$$

$$t \ll R \quad (R+t)^4 \approx R^4 + 4R^3 t$$

$$I = R^3 t \pi$$



Primer: enostavno podprti nosilec dolžine $l = 2\text{m}$ s tankostenskim krožnim presekom je točkovno obremenjen na svoji polovici. Določi debelino, da bo osna napetost pod dopustno vrednostjo.

$$\sigma \leq \sigma_0$$

$$\sigma_{\max} = \frac{M_{\max}}{I} z_{\max} < \sigma_0$$

$$I = \pi R^3 t$$

$$M_{\max} = \frac{1}{4} l F ; \quad z_{\max} = R + t$$



$$M_{\max} = \frac{1}{4} l F$$

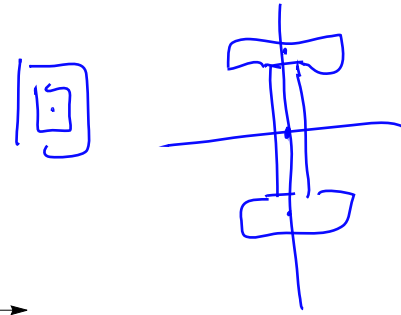
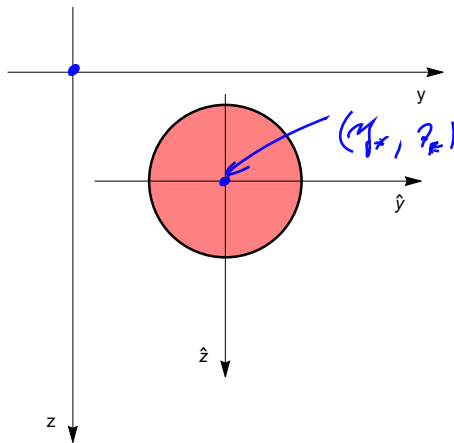
$$\frac{\frac{1}{4} l F \cdot (R+t)}{\pi R^3 t} < \sigma_0 \Rightarrow \frac{1}{4} l F (R+t) < \sigma_0 \pi R^3 t$$

$$\frac{1}{4} l F R < \sigma_0 \pi R^3 t - \frac{1}{4} l F t$$

$$\frac{1}{4} l F R < \pi (\sigma_0 R^3 t - \frac{1}{4} l F t)$$

Podpostavimo $\sigma_0 R^3 t > \frac{1}{4} l F$.

Potem mora veljati $t > \frac{l F R}{4 (\sigma_0 R^3 t - \frac{1}{4} l F)}$.



Slika 2: Paralelna koordinatna sistema.

$$y = y_* + \hat{y} \quad z = z_* + \hat{z} \quad dA = dy dz = d\hat{y} d\hat{z}$$

$$\int_A z^2 dA = \int_A (z_* + \hat{z})^2 dA = \int_A (z_*^2 + 2z_*\hat{z} + \hat{z}^2) dA =$$

$$= \underbrace{z_*^2}_{A} \int_A dA + 2z_* \int_A \hat{z} dA + \int_A \hat{z}^2 dA =$$

$$= z_*^2 A + \hat{I}$$

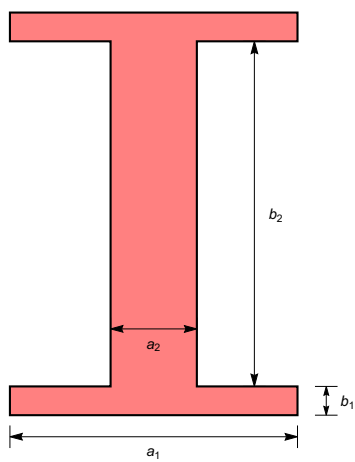
$$\hat{I} = \hat{I} + z_*^2 A$$

$$\int_A \hat{z} dA = \int_A \hat{z} d\hat{z} \int dy$$

||
0

Izrek o paralelnih oseh, $\hat{I} = z_0^2 A + I$.

Primer: ploskovni moment I nosilca.



Slika 3: I nosilec.